

(1) (a) A plane polarized EM wave with electric field given by $\vec{E} = \vec{e}_0 E_0 e^{i(kz - \omega t)}$ is incident normally on a semi-infinite slab of thickness d with its left surface at $z=0$ and dielectric constant $\epsilon(\omega)/\epsilon_0$ possessing finite conductivity.

With no assumptions regarding the smallness of $[\epsilon(\omega)/\epsilon_0 - 1]$, show that to *first order* in d , the electric field at point z downstream from the slab is given by

$$\vec{E} = \vec{e}_0 E_0 e^{i(kz - \omega t)} \left[1 + ik \left(\frac{\epsilon(\omega)}{\epsilon_0} - 1 \right) \frac{d}{2} \right]$$

(10 points)

(b) Calculate the rate of power dissipation in the slab per unit area to lowest order in the thickness d .

(15 points)

(2) Consider circularly polarized EM waves propagating in the direction of a static magnetic field \mathbf{B}_0 in a medium consisting of N electrons/unit volume behaving as bound oscillators with a single oscillator resonance frequency ω_0 with oscillator strength = 1 and damping constant γ .

(a) Show that the relationship between the refractive indices for waves of (+) and (-) circular polarization can be written as

$$n_+^2 - n_-^2 = \frac{Ne^2}{\epsilon_0 m} \times \left[\frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega + eB_0\omega / m} - \frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega - eB_0\omega / m} \right]$$

(12 points)

(b) Neglecting γ , from the above formula calculate the rotation in radians of the plane of polarization of a *linearly polarized* plane wave after propagating a length L in the medium when the magnetic field is present.

(Hint: Consider a linearly polarized wave as the coherent sum of two oppositely circularly polarized waves).

(13 points)

(3) The theory of magneto-optic effects shows that in a medium with magnetization density \mathbf{M} , the \mathbf{D} and \mathbf{E} fields are related via an anisotropic dielectric tensor $\epsilon_{\alpha\beta}$

$$D_\alpha = \sum_\beta \epsilon_{\alpha\beta} E_\beta$$

where $\epsilon_{\alpha\beta} = \epsilon_0 \delta_{\alpha\beta} + \delta\epsilon_{\alpha\beta}$

and the $\delta\epsilon_{\alpha\beta}$ are small and have the form:

$$\frac{\delta\epsilon_{\alpha\beta}}{\epsilon_0} = \begin{pmatrix} A & -iBM_z & iBM_y \\ iBM_z & A & -iBM_x \\ -iBM_y & iBM_x & A \end{pmatrix}$$

A and B are constants and M_x etc. are the Cartesian components of \mathbf{M} .

(a) Show that Maxwell's Equations for a wave propagating in the medium with frequency ω can be written as the set of equations

$$\nabla^2 E_\alpha + k^2 E_\alpha = \nabla_\alpha (\nabla \cdot \mathbf{E}) - k^2 \sum_\beta \frac{\delta\epsilon_{\alpha\beta}}{\epsilon_0} E_\beta$$

where $k = \omega/c$. (10 points)

(b) Using the outgoing Green's function formalism and the Born approximation (as in Secn. 10.2 in Jackson) show that the scattered electric field from a small volume of this material in the far field region can be written to first order in the $\delta\epsilon_{\alpha\beta}$ as

$$\mathbf{E} \rightarrow \mathbf{E}^0 + \mathbf{A}_{sc} \frac{e^{ikr}}{r}$$

$$\text{where } \frac{\mathbf{A}_{sc} \cdot \mathbf{e}_{sc}}{E_0} = \frac{k^2}{4\pi} \int d^3x e^{i\mathbf{q} \cdot \mathbf{x}} \sum_{\alpha\beta} (e_{sc})_\alpha \frac{\delta\epsilon_{\alpha\beta}}{\epsilon_0} (e_0)_\beta$$

where $\mathbf{q} = \mathbf{k}_0 - \mathbf{k}_{sc}$, (both \mathbf{k}_0 and \mathbf{k}_{sc} have magnitude k)

the spatial dependence of the incident electric field is given by $\mathbf{E}^0(\mathbf{x}) = \mathbf{e}_0 E_0 e^{i\mathbf{k}_0 \cdot \mathbf{x}}$

and \mathbf{e}_{sc} is one of the possible polarization vectors of the scattered beam. (10 points)

(c) Assume the incident wave is propagating along the z-direction with \mathbf{e}_0 along the x-direction and \mathbf{M} is parallel to the z-direction and the volume is a sphere of radius a . For scattering through an angle θ in the y-z plane, calculate $\frac{d\sigma}{d\Omega}$ explicitly for the outgoing electric field polarized in the y-z plane; and in the x-direction (i.e. perpendicular to the y-z plane). (15 points)

NOTE 1: Do NOT consider any contributions from $\delta\mu$ -type terms!

NOTE 2: The cross-sections cannot involve E_0

(4) Use the Kramers-Kronig relation to calculate the real part of $\epsilon(\omega)$, given the imaginary part of $\epsilon(\omega)$ for positive ω as

$$\text{Im}(\epsilon / \epsilon_0) = \lambda[\theta(\omega - \omega_1) - \theta(\omega - \omega_2)]$$

$$\omega_2 > \omega_1 > 0$$

$$\theta(x) = 0, x < 0; \theta(x) = 1, x > 0$$

Sketch the behavior of $\text{Re } \epsilon(\omega)$ and $\text{Im } \epsilon(\omega)$ as functions of ω .

(15 points)